MEASUREMENTS OF X_{max} AND TESTS OF HADRONIC INTERACTIONS WITH THE PIERRE AUGER OBSERVATORY

C. Bleve^{*} for the Pierre Auger Collaboration^{\dagger}

*Bergische Universität Wuppertal, FB C Physik, Gaußstr. 20, D-42097 Wuppertal, Germany
[†]Observatorio Pierre Auger, Av. San Martin Norte 304, 5613 Malargüe, Argentina
full author list at http://www.auger.org/archive/authors_2011_3.html

The Pierre Auger Observatory is designed to detect the particle showers produced in the atmosphere by the most energetic cosmic rays, particles with individual energies up to $E \simeq 10^{20}$ eV. The challenge is to measure their spectrum, arrival directions and mass composition. The depth of the maximum of the longitudinal shower development (X_{max}) is an indicator of the elemental composition. We present the measurement of the first two moments of the X_{max} distribution for $E > 10^{18}$ eV. The interpretation in terms of primary masses can only be done by comparison with predictions of hadronic interaction models. The description of hadronic interactions can be tested by determining the muon content of air shower data. Our results at $E \sim 10^{19}$ eV are compared with the predictions of the QGSjetII model for both proton and iron primaries, showing an observed excess of muons with respect to the model.

1 1 Introduction

Cosmic rays (CR) are a natural beam of ionized atomic nuclei with a rapidly falling energy 2 spectrum that extends up to $E \sim 10^{20}$ eV. To understand the sources and propagation of CRs, 3 the measurement of their flux, elemental composition and distribution of arrival directions is 4 needed. In the highest energy range of the CR spectrum, fluxes are too low for direct observation 5 with satellites or balloon-borne instruments: ground based detectors are used. They observe the 6 particle showers initiated by CRs when interacting with the terrestrial atmosphere. From the point of view of particle physics, the detection of CRs at extreme energies can be regarded as a fixed target collider experiment. The first interactions between primary CRs and atmospheric nuclei reach energies equivalent to p-p collisions at $\sqrt{s} \simeq 400$ TeV, about one order of magnitude 10 above those accessible with the LHC. 11

12 2 The Pierre Auger Observatory

The Pierre Auger Observatory operates in the Ultra High Energy (UHE) range ($E > 10^{18} eV$). 13 Its Southern Site is located in the Mendoza province, in Argentina. The Auger Observatory is 14 a hybrid experiment combining two complementary detection techniques. A $3,000 \text{ km}^2$ surface 15 detector (SD) samples the shower particles reaching ground level with 1660 water-Cherenkov 16 detectors. The SD is overlooked by a fluorescence detector (FD): 27 telescopes at 4 sites detect 17 the fluorescence light emitted along the longitudinal path of the shower. The SD has $\sim 100\%$ 18 duty cycle, and detection efficiency > 97% above 3 EeV for zenith angles $< 60^{\circ}$.² The signal in 19 each water-Cherenkov detector is recorded as FADC traces from 3 photomultipliers. The FD, 20



Figure 1: The measured X_{max} and $\text{RMS}(X_{\text{max}})$ as a function of the primary energy are compared with air shower simulations using different primary particles and hadronic interaction models. The $\text{RMS}(X_{\text{max}})$ distribution is obtained after subtracting in quadrature the detector resolution.⁶

which operates in clear moonless nights (duty cycle $\tilde{1}5\%$), can detect CRs down to an energy of ~ 10^{18} eV and is able to observe, within its field of view, the longitudinal profile of the shower. The coexistence of the two detectors combines the high exposure of the SD, yelding high statistics, and the almost calorimetric measurement of the energy of air showers provided by the FD. A subset of high quality hybrid events (golden hybrids), detected and reconstructed independently by the SD and FD, are used to calibrate the SD energy scale on the FD one.³ The estimate of the total systematic uncertainty on the energy scale is 22%.⁴

28 **3** Measurements of X_{max}

²⁹ The depth at which the longitudinal development of the shower reaches its maximum contains ³⁰ information about the mass of the primary CR initiating the shower and about the properties ³¹ of hadronic interactions. It can be measured with high accuracy by the FD. The average value ³² $\langle X_{\text{max}} \rangle$ depends on the primary energy E and on the number of nucleons A:

$$\langle X_{\max} \rangle = \alpha \left(\ln E - \langle \ln A \rangle \right) + \beta, \tag{1}$$

with α and β depending on the the details of hadronic interactions. Their values are very 33 sensitive to changes in cross-section, multiplicity and elasticity⁵ Eq. 1 is derived from a simple 34 generalization of the Heitler model to showers induced by hadronic primaries, but it provides a 35 good description of the $X_{\rm max}$ evolution predicted by hadronic models currently in use. In the 36 energy range of interest for the Auger Observatory, α and β can be considered independent of E. 37 Another mass sensitive quantity is $RMS(X_{max})$, expressing quantitatively the shower-to-shower 38 fluctuations of X_{max} . It is expected to decrease with the number of nucleons A. The measurement 39 of X_{max} and $\text{RMS}(X_{\text{max}})$ presented in Fig. 1 is based on the analysis of hybrid data collected 40 between December 2004 and March 2009⁶. Hybrids are defined as events observed by the FD and 41 at least one SD station. After quality cuts, 3754 hybrid events are used. The number of events 42 per energy bin is shown in Fig. 2, left. A comparison with four widely used high-energy hadronic 43 interaction models ^{7,8,9,10} suggests a gradual transition to heavier composition with increasing 44 primary energy (Fig. 1, left). In the simple hypothesis of two mass components, however, the 45 $RMS(X_{max})$ results from the RMS of individual species and from the separation of their $\langle X_{max} \rangle$: 46 a gradual transition form p to Fe primaries should lead to an increase of $RMS(X_{max})$ above the 47 value predicted for protons.¹¹ This effect is not observed in the Auger measurements (Fig. 1, 48 right). 49

The elongation rate, defined as the variation of X_{max} per decade of energy, is sensitive to changes in composition with energy. A constant elongation rate cannot describe the measured evolution of $\langle X_{\text{max}} \rangle$ with energy. A broken line is used in Fig. 2 (left) to fit the distribution: a change of

 82^{+25}_{-21} g/cm²/decade in the elongation rate occurs at $\log_{10}(E/eV) = 18.24 \pm 0.05$. This is close 53 to the energy of the spectral ankle³ at $\log_{10}(E_{ankle}/eV) = 18.65 \pm 0.09(stat)^{+0.10}_{-0.11}(sys)$, which is 54 usually interpreted in terms of transition from galactic to extragalactic CRs. This interpretation 55 is supported by the observed change in the elongation rate, under the assumption that hadronic 56 interactions are not significantly changing with energy. 57

Muon content of air showers 4 58

The sensitivity of the Auger Observatory to both the electromagnetic and muonic components 59 of air showers, allows us to test predictions of their relative contributions given by hadronic 60 interaction models. The detector cannot, at present, measure separately the muon content of 61 air showers, but several methods have been developed for an indirect estimate of the contribution: 62

(a) the **universality method**¹², based on the hypothesis that the electromagnetic and muonic 63 signals measured at the ground for a fixed energy and a given distance from the shower axis 64 have a universal dependence on the difference in grammage between the observation level 65 and the shower maximum. The relative normalization of the muonic component can then 66 be determined by requiring the total signal to match the experimental observations; 67

- (b) the **jump method**¹³ correlates the signal differences between consecutive bins of the SD 68 FADC traces and the muon fraction of the total signal; 69
- (c) the **smoothing method**¹⁴ finds muon-induced peaks in the FADC traces of the SD stations 70 through an iterative smoothing procedure; 71

(d) the **golden hybrid analysis**¹⁴ selects, in a set of simulations with the same energy and 72 geometry of a given event, the longitudinal profile that matches most closely the measured 73 one. The simulated SD response is then compared with the SD measurement. 74

Events selected in the energy range $\log_{10}(E/eV) = 19 \pm 0.2$ ($\sqrt{s} \simeq 140$ TeV for proton primaries) 75 and zenith $\theta < 50^{\circ}$ have been analysed to obtain the muonic content at 1000 m from the shower 76

axis.¹⁴ The number of muons N_{μ}^{rel} relative to the prediction of QGS jet II⁹ for protons is shown in 77

Fig. 2 (right). The results obtained with the different methods are compatible with each other. 78



Figure 2: Left: A broken line is used to fit the evolution of $\langle X_{\text{max}} \rangle$ with the logarithm of the primary energy. The change of slope (elongation rate) supports a change in the mass composition of cosmic rays at the ankle. The measurements of X_{max} published by the HiRes collaboration¹⁵ are shown for comparison. Right: The muon content of $E=10^{19\pm0.2}$ eV air showers is determined at 1000 m from the shower axis with different indirect methods (a-d, see text and references therein) and expressed as a number relative to the prediction of the QGSjetII model for protons primaries. That number is shown as a function of the shift in energy scale with respect to the FD

⁷⁹ Methods (a) and (c) constrain the energy scale to values higher than the FD scale but compatible ⁸⁰ with its systematic uncertainty. If the energy is fixed at this new scale, the muon content in the ⁸¹ data exceeds the QGSjetII prediction for protons by 30% to 70%. The model prediction for iron ⁸² primaries ($N_{\mu}^{rel} = 1.32$, blue horizontal line) is marginally compatible with the results of the ⁸³ methods within their systematic uncertainties. The measurement of X_{max} at the same energy, ⁸⁴ hovever, is not compatible with QGSjetII predictions for a pure iron composition (Sec. 3).

 $_{\tt 85}$ The AMIGA 16 extension of the Auger Observatory (in R&D) will employ scintillation counters

⁸⁶ buried 2.3 m underground to avoid the detection of the electromagnetic component of air show-⁸⁷ ers. AMIGA will provide, on a fraction of the SD area, a direct measurement of the muonic

component that can be used to test and calibrate the indirect methods used so far.

89 5 Conclusions

The Pierre Auger Observatory is the largest detector currently in operation for the detection of 90 cosmic rays in the UHE range. Its main scientific goal, from the point of view of astrophysics, is 91 to find clues about the sources, acceleration and propagation of cosmic rays. From the point of 92 view of particle physics, it observes, through air showers, collisions up to $\sqrt{s} \simeq 400$ TeV, where 93 the properties of particle interactions have large uncertainties and are extrapolated from collider 94 measurements using hadronic interaction models. On the one hand reliable models are required 95 for a precise interpretation of data, in particular for determining the mass composition of CRs, 96 on the other hand the data collected with the Auger Observatory provide a unique test bed 97 for constraining model predictions at extreme energies. The depth of the shower maximum, its 98 fluctuations and the muon content of showers at ground level are important observables related 99 to the elemental composition and are model dependent. The measurement of $X_{\rm max}$ suggests a 100 transition toward increasingly heavier primaries with increasing energies, although the measured 101 $RMS(X_{max})$ is smaller than models would predict. The QGS jet II model was tested against data: 102 at ~ 10^{19} eV it fails to describe consistently both $X_{\rm max}$ and muon content at 1000 m from the 103 axis. A deficit of muons is found in the predictions of the QGS jet II model. 104

105 **References**

- 106 1. The Pierre Auger Coll., *Nucl. Instrum. Methods* A **620**, 227 (2010).
- 2. The Pierre Auger Coll., Nucl. Instrum. Methods A 613, 29 (2010).
- ¹⁰⁸ 3. The Pierre Auger Coll., *Phys. Lett.* B **685**, 239 (2010).
- 4. C. Di Giulio [Pierre Auger Coll.], Proc. of 31st ICRC (2009) arXiv:0906.2189.
- ¹¹⁰ 5. R. Ulrich, *Phys. Rev.* D **83**, 054026 (2011).
- 6. The Pierre Auger Coll., *Phys. Rev. Lett.* **104**, 091101 (2010).
- 112 7. E. -J. Ahn *et al.*, *Phys. Rev.* D **80**, 094003 (2009).
- ¹¹³ 8. N.N. Kalmykov and S.S. Ostapchenko, Phys. Atom. Nucl. **56**, 346 (1993).
- ¹¹⁴ 9. S.S. Ostapchenko, Nucl. Phys. Proc. Suppl. **151**, 143 (2006).
- 115 10. T. Pierog and K. Werner, Phys. Rev. Lett. **101**, 171101 (2008).
- 116 11. M. Unger [Pierre Auger Coll.], Proc. of UHECR2010 (2010) arXiv:1103.5857.
- 117 12. F. Schmidt et al., Astropart. Phys. 29, 355 (2008).
- 118 13. M. Healy [Pierre Auger Coll.], Proc. of 30th ICRC (2007) vol. 4, p. 377.
- 119 14. A. Castellina [Pierre Auger Coll.], Proc. of 31st ICRC (2009), arXiv:0906.2319.
- 120 15. The HiRes Coll., *Phys. Rev. Lett.* **100**, 101101 (2008).
- 121 16. M. Platino [Pierre Auger Coll.], Proc. of 31st ICRC (2009) arXiv:0906.2354.